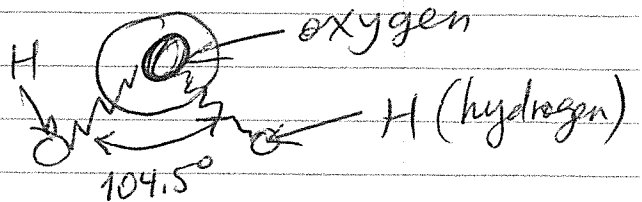


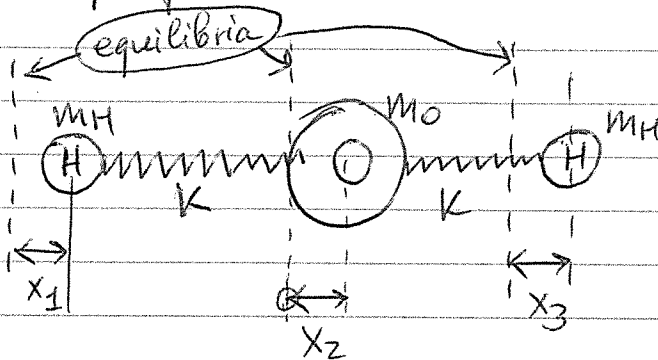
Last lecture in the course:

Application — oscillations of H_2O molecule

H_2O molecule:



Simplified view:



We approximate the restoring force between the O- and H-atoms using Hooke's law

$$F_{\text{spring}} = -k \cdot \Delta x \quad \leftarrow \text{change of spring's length.}$$

Use 2nd Law of Newton: $\mathbf{ma} = \mathbf{F}$

↑ acceleration,
 d^2x/dt^2 .

$$m_H \frac{d^2 x_1}{dt^2} = -k(x_1 - x_2)$$

$$m_O \frac{d^2 x_2}{dt^2} = -k(x_2 - x_1) - k(x_2 - x_3)$$

$$m_H \frac{d^2 x_3}{dt^2} = -k(x_3 - x_2)$$

The l.h.s. can be written as matrix-vector:

App-2

$$\underbrace{\begin{pmatrix} m_H & 0 \\ 0 & m_0 \\ 0 & m_H \end{pmatrix}}_{M, \text{ symmetric}} \underbrace{\frac{d^2}{dt^2} \begin{pmatrix} x_1 \\ x_2 \\ x_3 \end{pmatrix}}_{\underline{x}} = \underbrace{\begin{pmatrix} -k & k & 0 \\ k & -2k & k \\ 0 & k & -k \end{pmatrix}}_{K, \text{ symmetric}} \begin{pmatrix} x_1 \\ x_2 \\ x_3 \end{pmatrix} \quad (1a)$$

Can continue as:

$$\frac{d^2 \underline{x}}{dt^2} = (M^{-1} K) \underline{x}, \quad (1b)$$

but won't do that.

Reason: $(M^{-1} K)$ is not symmetric

(even though M and M^{-1} and K are symmetric, but $M^{-1} K$ is not because M^{-1} & K do not commute:

$A^T = A$, $B^T = B$, but $(AB)^T \neq AB$ since $AB \neq BA$; we did this in Sec. 1.6).

Although calculations can be done with a non-symmetric $(M^{-1} K)$, they are nicer with a symmetric matrix. Therefore, we proceed differently.

Note:
$$M = \begin{pmatrix} m_H & 0 \\ 0 & m_0 \\ 0 & m_H \end{pmatrix} = \begin{pmatrix} \sqrt{m_H} & 0 \\ 0 & \sqrt{m_0} \\ 0 & \sqrt{m_H} \end{pmatrix} \begin{pmatrix} \sqrt{m_H} & 0 \\ 0 & \sqrt{m_0} \\ 0 & \sqrt{m_H} \end{pmatrix}.$$

App-3

$$\text{So, } M = \sqrt{M} \cdot \sqrt{M}.$$

Note that \sqrt{M} is also symmetric:

$$\sqrt{M}^T = \sqrt{M} \quad (2)$$

Now, transform (1b) as follows:

$$\sqrt{M} \sqrt{M} \frac{d^2 \underline{x}}{dt^2} = K \sqrt{M}^{-1} \sqrt{M} \underline{x} \Rightarrow$$

$$\frac{d^2}{dt^2} \underbrace{(\sqrt{M} \underline{x})}_y = \underbrace{(\sqrt{M}^{-1} K \sqrt{M}^{-1})}_A (\sqrt{M} \underline{x}) \quad (3a)$$

$$\frac{d^2 y}{dt^2} = A y, \quad (3b)$$

where now A is symmetric. Indeed:

$$A^T = (\sqrt{M}^{-1} K \sqrt{M}^{-1})^T = (\sqrt{M}^{-1})^T K^T (\sqrt{M}^{-1})^T$$

\sqrt{M}^{-1} and K are symmetric $\Rightarrow \sqrt{M}^{-1} K \sqrt{M}^{-1} = A. \quad \checkmark$

Now let's find its explicit form:

$$A = \begin{pmatrix} \frac{1}{\sqrt{m_H}} & 0 \\ 0 & \frac{1}{\sqrt{m_0}} \\ 0 & \frac{1}{\sqrt{m_H}} \end{pmatrix} \begin{pmatrix} -K & K & 0 \\ K & -2K & K \\ 0 & K & -K \end{pmatrix} \begin{pmatrix} \frac{1}{\sqrt{m_H}} & 0 \\ 0 & \frac{1}{\sqrt{m_0}} \\ 0 & \frac{1}{\sqrt{m_H}} \end{pmatrix} =$$

$$= \begin{pmatrix} -\frac{\kappa}{m_H} & \frac{\kappa}{\sqrt{m_H m_0}} & 0 \\ \frac{\kappa}{\sqrt{m_H m_0}} & -\frac{2\kappa}{m_0} & \frac{\kappa}{\sqrt{m_H m_0}} \\ 0 & \frac{\kappa}{\sqrt{m_H m_0}} & -\frac{\kappa}{m_H} \end{pmatrix} \stackrel{\text{notation}}{\equiv} \underbrace{\begin{pmatrix} -a & b & 0 \\ b & -2c & b \\ 0 & b & -a \end{pmatrix}}_A$$

(App-4)

Note:

$$ac = b^2 \quad (4)$$

$$\left(\frac{\kappa}{m_H} \cdot \frac{\kappa}{m_0} = \left(\frac{\kappa}{\sqrt{m_H m_0}} \right)^2 \right)$$

We now seek solution of (3b) in the form:

$$\underline{y}(t) = e^{\sqrt{\lambda}t} \underline{u} \leftarrow \text{independent of } t. \quad (5)$$

(App-3)

Substitute (5) into (3b):

$$\cancel{(\sqrt{\lambda})^2 e^{\sqrt{\lambda}t}} \underline{u} = A \cancel{e^{\sqrt{\lambda}t}} \underline{u} \Rightarrow$$

$$A \underline{u} = \lambda \underline{u} \quad (6)$$

The eigenvalue problem!

Its eigenvalues ^{and eigenvectors} are found ~~by~~ as usual (with some effort, but straightforwardly):

App-5

$$\lambda_1 = -a, \quad \underline{u}_1 = \begin{pmatrix} -1 \\ 0 \\ 1 \end{pmatrix}$$

$$\lambda_2 = 0, \quad \underline{u}_2 = \begin{pmatrix} 1/b \\ 1 \\ 1 \end{pmatrix} = \begin{pmatrix} 1/\sqrt{m_0/m_H} \\ 1 \\ 1 \end{pmatrix} \quad (7)$$

$$\lambda_3 = -(2c+a), \quad \underline{u}_3 = \begin{pmatrix} 1 \\ -2b/a \\ 1 \end{pmatrix} = \begin{pmatrix} 1 \\ -2\sqrt{m_H/m_0} \\ 1 \end{pmatrix}$$

As the general theory predicts, they are all orthogonal.

Returning to Eq. (5) on p. App-4, we see that we have found 3 solutions of the diff. equation (36) (p. App-3):

$$y_1 = e^{\sqrt{\lambda_1}t} \underline{u}_1, \quad y_2 = e^{\sqrt{\lambda_2}t} \underline{u}_2, \quad y_3 = e^{\sqrt{\lambda_3}t} \underline{u}_3.$$

How do we use them to construct the general solution of the diff. eq. and also its solution with a particular initial condition?

We use the fact that the diff. eq. (36) is linear and homogeneous. Then its solutions form a linear vector space (Sec. 3.2). This means that any linear combination of solutions is also a solution of (36)!

App-6

So the most general solution of (3b) is:

$$y(t) = c_1 e^{\sqrt{\lambda_1} t} \underline{u}_1 + c_2 e^{\sqrt{\lambda_2} t} \underline{u}_2 + c_3 e^{\sqrt{\lambda_3} t} \underline{u}_3$$

(simplification: we dropped terms $e^{-\sqrt{\lambda_i} t} \underline{u}_i$ etc.) (8)

How do we find ~~the~~ c_1, c_2, c_3 ?
From the initial condition!

$$y(0) = c_1 \cdot e^{\sqrt{\lambda_1} \cdot 0} \underline{u}_1 + c_2 \underline{u}_2 + c_3 \underline{u}_3 \quad (9)$$

↑ known i.c.

~~the~~ (Note: This is a simplification, but the basic idea is correct.) (see note after 8)

Since $\underline{u}_1, \underline{u}_2, \underline{u}_3$ are orthogonal, the coordinates c_1, c_2, c_3 are found easily:

$$c_j = \frac{\underline{u}_j^T y(0)}{\underline{u}_j^T \underline{u}_j} \quad (10)$$

Thus, $y(t) = \sqrt{M} x(t)$ has been found.

Let us now write the answer in terms of x and interpret it.

App-7

$$\underline{y} = \sqrt{M} \underline{x} \Rightarrow \underline{u} = \sqrt{M} \underline{v} \rightarrow \begin{pmatrix} x_1 \\ x_2 \\ x_3 \end{pmatrix}$$

$$\Rightarrow \underline{v}_j = \sqrt{M}^{-1} \underline{u}_j$$

$$\textcircled{1} \underline{u}_1 = \begin{pmatrix} -1 \\ 0 \\ 1 \end{pmatrix} \Rightarrow \underline{v}_1 = \begin{pmatrix} 1/\sqrt{m_H} & 0 & 0 \\ 0 & 1/\sqrt{m_0} & 0 \\ 0 & 0 & 1/\sqrt{m_H} \end{pmatrix} \begin{pmatrix} -1 \\ 0 \\ 1 \end{pmatrix} \Rightarrow$$

$$\Rightarrow \underline{v}_1 = \begin{pmatrix} -1/\sqrt{m_H} \\ 0 \\ 1/\sqrt{m_H} \end{pmatrix} \sim \begin{pmatrix} -1 \\ 0 \\ 1 \end{pmatrix}$$

$$\textcircled{2} \underline{u}_2 = \begin{pmatrix} 1 \\ \sqrt{m_0/m_H} \\ 1 \end{pmatrix} \Rightarrow \underline{v}_2 = \begin{pmatrix} 1/\sqrt{m_H} & 0 & 0 \\ 0 & 1/\sqrt{m_0} & 0 \\ 0 & 0 & 1/\sqrt{m_H} \end{pmatrix} \begin{pmatrix} 1 \\ \sqrt{m_0/m_H} \\ 1 \end{pmatrix}$$

$$\underline{v}_2 = \begin{pmatrix} 1/\sqrt{m_H} \\ 1/\sqrt{m_H} \\ 1/\sqrt{m_H} \end{pmatrix} \sim \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix}$$

$$\textcircled{3} \underline{u}_3 = \begin{pmatrix} 1 \\ -2\sqrt{m_H/m_0} \\ 1 \end{pmatrix} \Rightarrow \underline{v}_3 = \dots \begin{pmatrix} 1/\sqrt{m_H} \\ -2\sqrt{m_H/m_0} \\ 1/\sqrt{m_H} \end{pmatrix} \sim \begin{pmatrix} 1 \\ -2\frac{m_H}{m_0} \\ 1 \end{pmatrix}$$

$$m_0/m_H \approx 18/1, \Rightarrow \underline{v}_3 = \begin{pmatrix} 1 \\ -1/8 \\ 1 \end{pmatrix}$$

So

$$\underline{x}(t) = c_1 e^{\sqrt{-a}t} \begin{pmatrix} -1 \\ 0 \\ 1 \end{pmatrix} + c_2 e^{at} \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix} + c_3 e^{\sqrt{-2a}t} \begin{pmatrix} 1 \\ -1/8 \\ 1 \end{pmatrix}$$

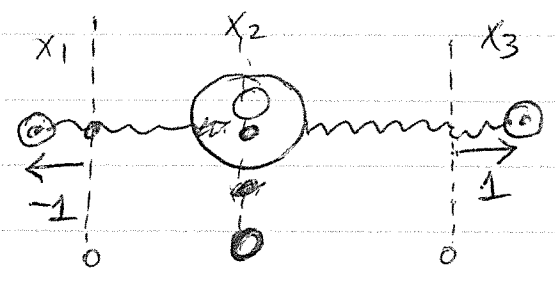
$$\sqrt{-\frac{k}{m_H}} = i\sqrt{\frac{k}{m_H}}$$

(11)

$$e^{i\alpha t} = \cos(\alpha t) + i \sin(\alpha t) \leftarrow \text{oscillations}$$

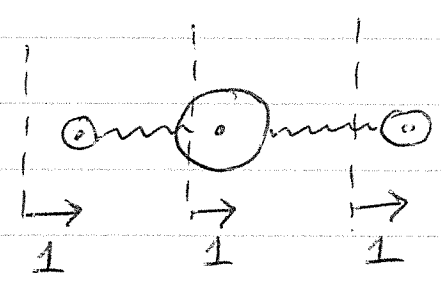
Interpretation:

$$\underline{v}_1 = \begin{pmatrix} -1 \\ 0 \\ 1 \end{pmatrix}$$



Center remains fixed.

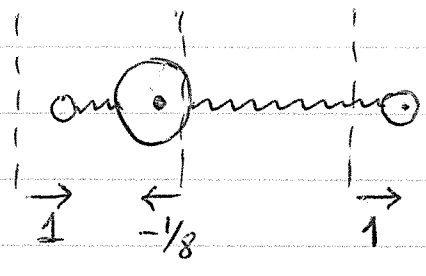
$$\underline{v}_2 = \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix}$$



This is simply a translation!

This is why its time dependence is $\sim e^{0 \cdot t}$ (time independent!) — no oscillations.

$$\underline{v}_3 = \begin{pmatrix} 1 \\ -1/8 \\ 1 \end{pmatrix}$$



Asymmetric oscillation.

Thus: The motion of a 3-atomic molecule is a linear combination of 2 oscillatory modes ($\underline{v}_1, \underline{v}_3$) and 1 translation (\underline{v}_2).