

## Sec. 16.9. Divergence Thm.

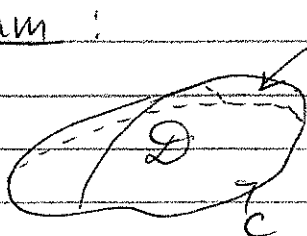
(35-1)

(Gauss Thm.)

① Main result.

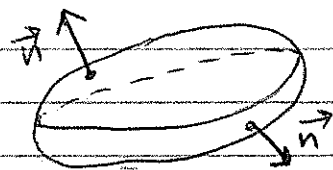
It extends Green's / Stokes Thm. to 3D.

Stokes Thm:



$$\oint_C \vec{F} \cdot d\vec{r} = \iint_S \text{curl } \vec{F} \cdot d\vec{S}$$

Divergence Thm:



$E = 3D \text{ solid}$

$S = \text{its boundary}$   
(closed surface)

"O" stresses that  $S$  is closed.

$$\oint_S \vec{F} \cdot d\vec{S} = \iiint_E \text{div } \vec{F} \, dV$$

Flux through  $S$ .

Thm: Let  $E$  be a simply connected solid and  $S$  be the boundary (surface) of  $E$ , with positive orientation (outside). Let  $\vec{F}$  have continuous partial derivatives in an open region that contains  $E$ .

Then:

$$\oint_S \vec{F} \cdot d\vec{S} = \iiint_E \text{div } \vec{F} \, dV$$

Flux through boundary of  $E$ .

Make sure to write

$\oint$ ,

not just  $\iint$ .

Note: This is an extension of FTC to 3D; see Sec. 16.4; also the Summary (Sec 16.10) in textbook

Proof

0) Rewrite the statement.

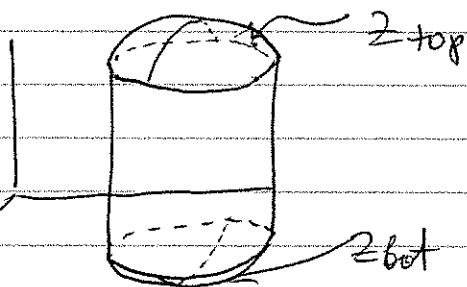
$$\vec{F} = \langle P, Q, R \rangle; \quad d\vec{S} \equiv \vec{n} dS.$$

Then

$$\iiint_E \left( \frac{\partial P}{\partial x} + \frac{\partial Q}{\partial y} + \frac{\partial R}{\partial z} \right) dV = \iint_S (P\vec{i} + Q\vec{j} + R\vec{k}) \cdot \vec{n} dS.$$

Let's prove the  $R$ -part; the rest are similar.

1) Suppose  $E$  is bounded by top, bottom, and vertical "walls"

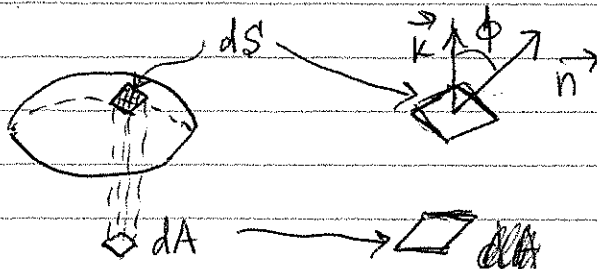


$$\iiint_E \frac{\partial R}{\partial z} dV = \iint_D \left( \int_{z_{\text{bot}}}^{z_{\text{top}}} \frac{\partial R}{\partial z} dz \right) dx dy$$

↑ projection of  $E$  on  $xy$ -plane

$$= \iint_D (R(x, y, z_{\text{top}}) - R(x, y, z_{\text{bot}})) dx dy.$$

2)



$$dA = dS \cdot \cos \phi$$

(see Sec. 16.7)

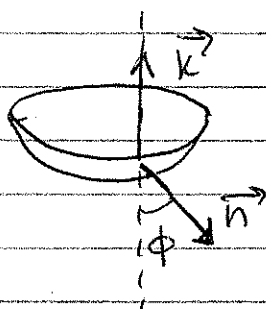
Then at the top surface:

$$\iint_{\mathcal{D}} R(x, y, z_{\text{top}}) dA = \iint_{\mathcal{D}} R(x, y, z_{\text{top}}) \cdot \cos \phi \cdot dS$$

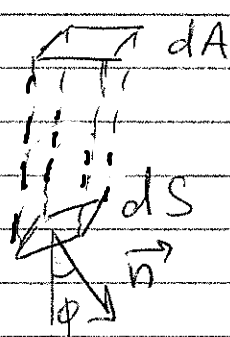
$$= \iint_{S_{\text{top}}} R(x, y, z) \vec{k} \cdot \vec{n} \, dS.$$

$\uparrow$   
 on  $S_{\text{top}}, z = z_{\text{top}}!$

3) At the bottom surface:



Now, again,  $dA = dS \cdot \cos \phi$ ,



but

$$\cos \phi = (-\vec{k} \cdot \vec{n}),$$

because  $\vec{k} \cdot \vec{n} < 0$ .

$$\text{So } \iint_{\mathcal{D}} R(x, y, z_{\text{bot}}) dA = \iint_{S_{\text{bot}}} R(x, y, z) (-\vec{k} \cdot \vec{n}) dS$$

4) Thus,

$$\iiint_E \frac{\partial R}{\partial z} dV = \iint_{\mathcal{D}} (R(x, y, z_{\text{top}}) - R(x, y, z_{\text{bot}})) dA$$

$$= \iint_{S_{\text{top}}} R \vec{k} \cdot \vec{n} \, dS - \iint_{S_{\text{bot}}} R \cdot (-\vec{k} \cdot \vec{n}) \, dS$$

$$= \iint_{(S_{\text{top}} + S_{\text{bot}})} R \cdot \vec{k} \cdot \vec{n} \, dS \quad \left( \begin{array}{l} \text{on the side walls, } \vec{n} \perp \vec{k} \\ \Rightarrow \vec{k} \cdot \vec{n} = 0 \end{array} \right)$$

$\Rightarrow$

$$\iiint_E \frac{\partial R}{\partial z} dV = \oint_{S' = \text{full boundary of } E} R \vec{k} \cdot \vec{n} dS.$$

### Significance of Div. theorem:

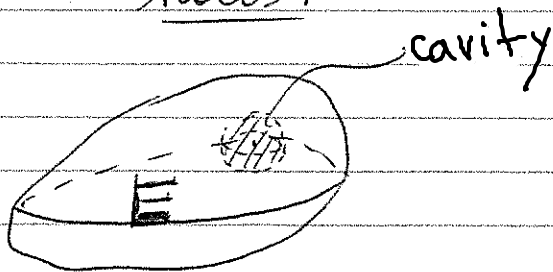
Can compute the flux of  $\vec{F}$  through closed surface  $S$  by evaluating the

$$\iiint_E \text{div} \vec{F} dV, \text{ where } E \text{ is bounded by } S.$$

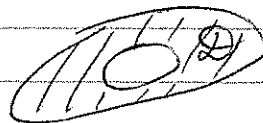
(Often,  $\text{div} \vec{F}$  has a simpler form than  $\vec{F} \cdot \vec{n}$ .)

See Ex. 1, 2 in book.

### ② Extension of Div. Thm. when $E$ has holes.



Similarly to a derivation in Sec. 16.4 where  $\mathcal{Q}$  had a hole:



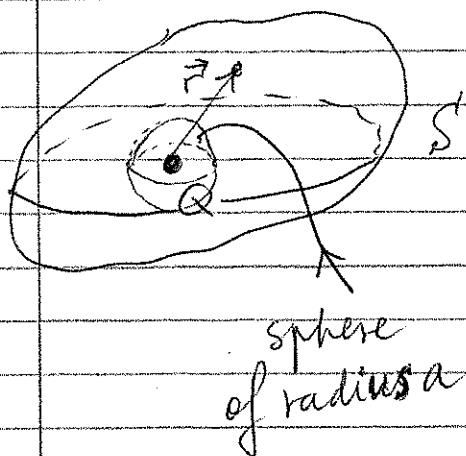
here we have:

$$\iiint_E \text{div} \vec{F} dV = \left( \oint_{S_{\text{outside}}} + \oint_{S_{\text{cavity}}} \right) \vec{F} \cdot \vec{n} dS,$$

where  $\vec{n}$  points outside on  $S_{\text{outside}}$ ;  
 $\vec{n}$  points inside on  $S_{\text{cavity}}$ .

Ex. 1

In electrostatics,  
the following problem is  
important.



Find the flux of electric  
field  $\vec{E}$  through any  
surface  $S$  enclosing a  
point charge  $Q$ .

Sol'n: 1)  $\vec{E} = \frac{Q\vec{r}}{|\vec{r}|^3}$ ,  $\vec{r} = \langle x, y, z \rangle$ .

2)  $\vec{E}$  (and  $\text{div } \vec{E}$ ) are not continuous at  
 $\vec{r} = 0$ ,  $\Rightarrow$  cannot apply the Div. Thm.  
to the region enclosed by  $S$ .

But can apply it to a region with a  
hole around  $\vec{r} = 0$ !

So:  $\left( \oint_S + \oint_{S_{\text{cavity}}} \right) \vec{E} \cdot d\vec{S} = \iiint_{\substack{\vec{E} \text{ with a cavity}}} \text{div } \vec{E} dV$

3) Verify by direct calculation (tedious!)  
that

$$\text{div } \vec{E} = Q \text{div } \frac{\langle x, y, z \rangle}{\sqrt{x^2 + y^2 + z^2}^3} = 0 \quad (\text{for } \vec{r} \neq 0)$$

Then  $\left( \oint_S + \oint_{S_{\text{cavity}}} \vec{E} \cdot d\vec{S} \right) = \iiint 0 dV = 0$ .

So  $\oint_S \vec{E} \cdot d\vec{S} = - \oint_{S_{\text{cavity}}} \vec{E} \cdot d\vec{S} = + \oint_{S_{\text{cavity}}} \vec{E} \cdot d\vec{S}$

surface of the cavity  
with  $\vec{n}$  pointing outside

4) On  $S$ ,  $|\vec{r}| = a$ ,  $\vec{n} = \frac{\vec{r}}{|\vec{r}|}$ . Then

$$\begin{aligned} \oint_{S_{\text{cavity}}} \vec{E} \cdot d\vec{S} &= \oint_{S_{\text{cavity}}} \frac{Q\vec{r}}{|\vec{r}|^3} \cdot \frac{\vec{r}}{|\vec{r}|} dS \\ &= Q \oint_{S_{\text{cav.}}} \frac{\vec{r} \cdot \vec{r}}{|\vec{r}|^4} dS = Q \oint_{S_{\text{cav.}}} \frac{1}{a^2} dS \\ &= \frac{Q}{a^2} \oint_{S_{\text{cavity}}} dS = \frac{Q}{a^2} \cdot 4\pi a^2 = 4\pi Q. \end{aligned}$$

surface of a sphere of radius  $a$

Thus:

$$\oint_S \frac{Q\vec{r}}{|\vec{r}|^3} \cdot d\vec{S} = 4\pi Q$$

for any surface surrounding  $Q$

GAUSS  
Law

closed